

# The Calderón problem with partial data

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# Calderón problem

Medical imaging, Electrical Impedance Tomography:

$$\begin{cases} \operatorname{div}(\gamma(x)\nabla u) = 0 & \text{in } \Omega, \\ u = f & \text{on } \partial\Omega \end{cases}$$

where  $\Omega \subset \mathbb{R}^n$  bounded domain,  $\gamma \in L^\infty(\Omega)$  positive.

Boundary measurements given by DN map

$$\Lambda_\gamma : f \mapsto \gamma \partial_\nu u|_{\partial\Omega}.$$

**Inverse problem:** given  $\Lambda_\gamma$ , determine  $\gamma$ .

# Calderón problem

Model case of inverse boundary problems for elliptic equations (Schrödinger, Maxwell, elasticity).

Related to:

- ▶ optical tomography
- ▶ inverse scattering
- ▶ travel time tomography and boundary rigidity
- ▶ hybrid imaging methods
- ▶ invisibility

# Calderón problem

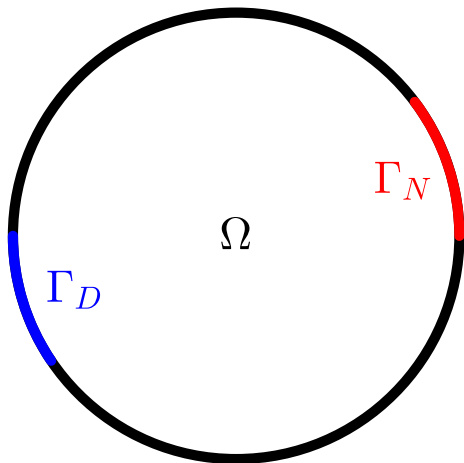
Uniqueness results:

- ▶ Calderón (1980): linearized problem
- ▶ Sylvester-Uhlmann (1987):  $n \geq 3$ ,  $\gamma \in C^2(\overline{\Omega})$
- ▶ Nachman (1996):  $n = 2$ ,  $\gamma \in W^{2,p}(\Omega)$
- ▶ Astala-Päivärinta (2006):  $n = 2$ ,  $\gamma \in L^\infty(\Omega)$
- ▶ Haberman-Tataru (2012):  $n \geq 3$ ,  $\gamma \in C^1(\overline{\Omega})$

We are interested in the *partial data problem* where measurements are made only on subsets of the boundary.

# Partial data problem

Prescribe voltages on  $\Gamma_D$ , measure currents on  $\Gamma_N$ :



# Partial data problem

Let  $\Gamma_D$  and  $\Gamma_N$  be open subsets of  $\partial\Omega$ . Define *partial Cauchy data set*

$$C_{\gamma}^{\Gamma_D, \Gamma_N} = \{(u|_{\Gamma_D}, \gamma \partial_{\nu} u|_{\Gamma_N}); \operatorname{div}(\gamma \nabla u) = 0 \text{ in } \Omega, u \in H^1(\Omega), \operatorname{supp}(u|_{\partial\Omega}) \subset \Gamma_D\}.$$

Corresponds to prescribing Dirichlet data on  $\Gamma_D$  and measuring Neumann data on  $\Gamma_N$ .

**Inverse problem:** given  $C_{\gamma}^{\Gamma_D, \Gamma_N}$ , determine  $\gamma$ .

# Partial data problem

Substitution  $u = \gamma^{-1/2}v$  reduces conductivity equation  $\operatorname{div}(\gamma \nabla u) = 0$  to *Schrödinger equation*  $(-\Delta + q)v = 0$ .

If  $q \in L^\infty(\Omega)$ , define

$$C_q^{\Gamma_D, \Gamma_N} = \{(u|_{\Gamma_D}, \partial_\nu u|_{\Gamma_N}); (-\Delta + q)u = 0 \text{ in } \Omega, u \in H_\Delta(\Omega), \operatorname{supp}(u|_{\partial\Omega}) \subset \Gamma_D\}.$$

Here  $H_\Delta(\Omega) = \{u \in L^2(\Omega); \Delta u \in L^2(\Omega)\}$ .

**Inverse problem:** given  $C_q^{\Gamma_D, \Gamma_N}$ , determine  $q$ .

# Partial data problem

Four main approaches for uniqueness:

1. Carleman estimates (Kenig-Sjöstrand-Uhlmann 2007)
2. Reflection approach (Isakov 2007)
3. 2D case (Imanuvilov-Uhlmann-Yamamoto 2010)
4. Linearized case (Dos Santos-Kenig-Sjöstrand-Uhlmann 2009)

The first two approaches work in dimensions  $n \geq 3$ . Will describe them in more detail.



# Strategy of proof

## Lemma (Integration by parts)

If  $\Gamma_D, \Gamma_N \subset \partial\Omega$  are open and if  $C_{q_1}^{\Gamma_D, \Gamma_N} = C_{q_2}^{\Gamma_D, \Gamma_N}$ , then

$$\int_{\Omega} (q_1 - q_2) u_1 u_2 \, dx = 0$$

for any  $u_j$  satisfying  $(-\Delta + q_j)u_j = 0$  in  $\Omega$  and

$$\text{supp}(u_1|_{\partial\Omega}) \subset \Gamma_D, \quad \text{supp}(u_2|_{\partial\Omega}) \subset \Gamma_N. \quad (*)$$

To show  $q_1 = q_2$ , enough that the set of products of solutions

$$\{u_1 u_2; (-\Delta + q_j)u_j = 0 \text{ in } \Omega, \quad u_j \text{ satisfy } (*)\}$$

is dense in  $L^1(\Omega)$ .

# Strategy of proof

Use special *complex geometrical optics solutions*

$$u \approx e^{\pm\tau\varphi} a, \quad (-\Delta + q)u = 0, \quad \text{supp}(u|_{\partial\Omega}) \subset \Gamma_{D,N}.$$

Here  $\tau > 0$  is a large parameter and

$$\left\{ \lim_{\tau \rightarrow \infty} u_1 u_2 \right\} \text{ dense in } L^1(\Omega).$$

Here  $\varphi$  is a *limiting Carleman weight*: Carleman estimate

$$\|e^{\pm\tau\varphi} v\|_{L^2(\Omega)} \leq \frac{C}{\tau} \|e^{\pm\tau\varphi} (-\Delta + q)v\|_{L^2(\Omega)}, \quad v \in C_c^\infty(\Omega).$$

(Also need boundary terms.) The function  $a$  is an *amplitude*.

# Strategy of proof

Condition for a *limiting Carleman weight*  $\varphi$ ,  $\nabla\varphi \neq 0$ :

$$\|e^{\pm\tau\varphi} v\|_{L^2(\Omega)} \leq \frac{C}{\tau} \|e^{\pm\tau\varphi} \Delta v\|_{L^2(\Omega)}, \quad v \in C_c^\infty(\Omega), \quad \tau \gg 1.$$

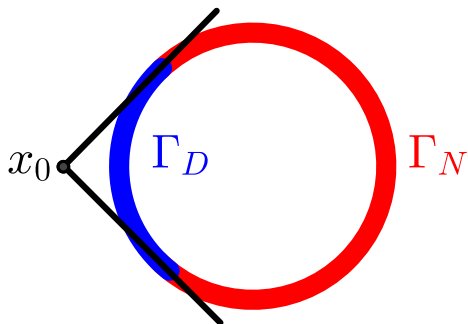
Results from Dos Santos-Kenig-S-Uhlmann (2009):

- ▶ conformally invariant condition
- ▶ if  $n \geq 3$ , only six basic forms for  $\varphi$ :

$$x_1, \quad \log|x|, \quad \frac{x_1}{|x|^2}, \quad \arctan \frac{x_2}{x_1}, \quad \dots$$

- ▶ if  $n = 2$ , any harmonic function is OK

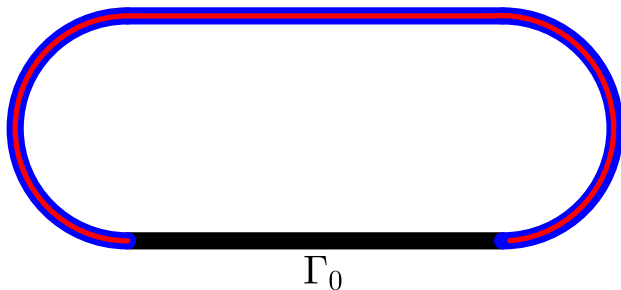
# Carleman estimate approach (KSU 2007)



- ▶  $\Gamma_D$  and  $\Gamma_N$  roughly complementary, need to overlap
- ▶  $\Gamma_D$  can be very small, but then  $\Gamma_N$  has to be very large
- ▶ proof uses weights  $\varphi(x) = \log|x - x_0|$  and Carleman estimates with boundary terms

## Reflection approach (Isakov 2007)

$$\Gamma_D = \Gamma_N = \Gamma$$



- ▶ local data:  $\Gamma_D = \Gamma_N = \Gamma$ , no measurements needed on  $\Gamma_0$
- ▶ the *inaccessible part of the boundary*,  $\Gamma_0$ , has strict restrictions (part of a hyperplane or part of a sphere)
- ▶ proof uses weights  $\varphi(x) = x_1$  and reflection about  $\Gamma_0$

## 2D and linearized cases

### Theorem (Imanuvilov-Uhlmann-Yamamoto 2010)

Let  $\Omega \subset \mathbb{R}^2$  be a bounded domain and let  $\Gamma \subset \partial\Omega$  be open. If  $q_1, q_2 \in C^{4,\alpha}(\overline{\Omega})$  for some  $\alpha > 0$  and if

$$C_{q_1}^{\Gamma,\Gamma} = C_{q_2}^{\Gamma,\Gamma},$$

then  $q_1 = q_2$ .

### Theorem (Dos Santos-Kenig-Sjöstrand-Uhlmann 2009)

Let  $\Omega \subset \mathbb{R}^n$  be a bounded domain,  $n \geq 2$ , and let  $\Gamma \subset \partial\Omega$  be open. The Cauchy data set  $C_q^{\Gamma,\Gamma}$  linearized at  $q = 0$  uniquely determines  $q$ .

# New results

Recall main approaches:

1. Carleman estimates
2. Reflection approach
3. 2D case
4. Linearized case

We unify approaches 1 and 2 and extend both. In particular, we relax the requirements on the inaccessible part in 2, and allow to use complementary (sometimes disjoint) sets as in 1.

The methods work for  $n \geq 3$ , also on certain Riemannian manifolds, and sometimes reduce the question to integral geometry problems of independent interest.

# New results

The first results are local results: given measurements on  $\Gamma \subset \partial\Omega$ , coefficients are determined in a neighborhood of  $\Gamma$ .

Proof reduces to an integral geometry problem (*Helgason support theorem*): recover a function locally from its integrals over *lines*, *great circles*, or *hyperbolic geodesics* in a certain neighborhood.

Instead of being completely flat or spherical, the inaccessible part  $\Gamma_0$  can be *conformally flat only in one direction*, e.g.

- ▶ cylindrical set (leads to integrals over lines)
- ▶ conical set (integrals over great circle segments)
- ▶ surface of revolution (integrals over hyperbolic geodesics).



# Cylindrical sets

## Theorem (Kenig-S 2012)

Let  $\Omega \subset \mathbb{R} \times \Omega_0$  where  $\Omega_0 \subset \mathbb{R}^2$  is convex, let  $\Gamma = \partial\Omega \setminus \Gamma_0$ , and suppose that  $\Gamma_0$  satisfies

$$\Gamma_0 \subset \mathbb{R} \times (\partial\Omega_0 \setminus E)$$

for some open set  $E \subset \partial\Omega_0$ . If  $q_1, q_2 \in C(\overline{\Omega})$  and if

$$C_{q_1}^{\Gamma, \Gamma} = C_{q_2}^{\Gamma, \Gamma},$$

then  $q_1 = q_2$  in  $\overline{\Omega} \cap (\mathbb{R} \times \text{ch}_{\mathbb{R}^2}(E))$ .

Corresponds to  $\varphi(x) = x_1$ . Similar result obtained independently by Imanuvilov and Yamamoto (2012).

# Conical sets

## Theorem (Kenig-S 2012)

Let  $\Omega \subset \{r\omega; r > 0, \omega \in M_0\}$  where  $M_0 \subset S^2$  is convex, let  $\Gamma = \partial\Omega \setminus \Gamma_0$ , and suppose that  $\Gamma_0$  satisfies

$$\Gamma_0 \subset \{r\omega; r > 0, \omega \in \partial M_0 \setminus E\}$$

for some open set  $E \subset \partial M_0$ . If  $q_1, q_2 \in C(\overline{\Omega})$  and if

$$C_{q_1}^{\Gamma, \Gamma} = C_{q_2}^{\Gamma, \Gamma},$$

then  $q_1 = q_2$  in  $\overline{\Omega} \cap (\mathbb{R} \times \text{ch}_{S^2}(E))$ .

Corresponds to  $\varphi(x) = \log|x|$ . Convex hull in  $S^2$  taken with respect to great circle segments.

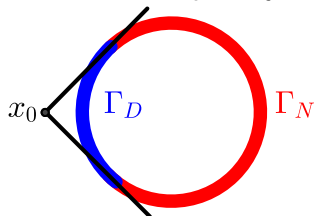
## Remarks

- ▶ convexity not required: if the inaccessible part is concave, recover the coefficient everywhere
- ▶ it is not required that  $\Gamma_D = \Gamma_N$ , can use somewhat complementary sets as in Kenig-Sjöstrand-Uhlmann
- ▶ sometimes  $\Gamma_D$  and  $\Gamma_N$  can be disjoint, for instance if

$$\Gamma_D = \{x \in \partial\Omega; (x - x_0) \cdot \nu(x) < 0\}$$

$$\Gamma_N = \{x \in \partial\Omega; (x - x_0) \cdot \nu(x) > 0\}$$

and if  $\{x \in \partial\Omega; (x - x_0) \cdot \nu(x) = 0\}$  has measure zero in  $\partial\Omega$ , then  $C_q^{\Gamma_D, \Gamma_N}$  determines  $q$  everywhere.

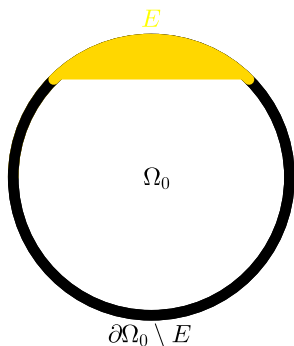


## Beyond the convex hull

Let  $\Omega \subset \mathbb{R} \times \Omega_0$  where  $\Omega_0 \subset \mathbb{R}^2$  is convex, let  $\Gamma = \partial\Omega \setminus \Gamma_0$ , and suppose that  $\Gamma_0$  satisfies

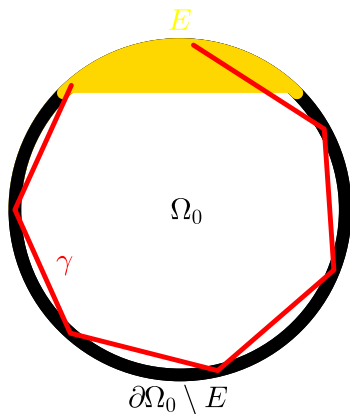
$$\Gamma_0 \subset \mathbb{R} \times (\partial\Omega_0 \setminus E)$$

for some open set  $E \subset \partial\Omega_0$ . From measurements on  $\Gamma$ , recover coefficient in  $\overline{\Omega} \cap (\mathbb{R} \times \text{ch}_{\mathbb{R}^2}(E))$ . Can one go beyond the convex hull?



## Beyond the convex hull

A continuous curve  $\gamma : [0, L] \rightarrow \overline{\Omega}_0$  is a *broken ray* if it consists of straight line segments that are reflected according to geometrical optics (angle of incidence = angle of reflection) when they hit  $\partial\Omega_0$ .



# Beyond the convex hull

## Theorem (Kenig-S 2012)

Let  $\Omega \subset \mathbb{R} \times \Omega_0$  where  $\Omega_0 \subset \mathbb{R}^2$  is a bounded domain, let  $\Gamma = \partial\Omega \setminus \Gamma_0$  where  $\Gamma_0$  satisfies for some open  $E \subset \partial\Omega_0$

$$\Gamma_0 \subset \mathbb{R} \times (\partial\Omega_0 \setminus E).$$

If  $q_1, q_2 \in C(\overline{\Omega})$  and  $C_{q_1}^{\Gamma, \Gamma} = C_{q_2}^{\Gamma, \Gamma}$ , then *for any nontangential broken ray*  $\gamma : [0, L] \rightarrow \overline{\Omega}_0$  *with endpoints on*  $E$ , and given any real number  $\lambda$ , one has

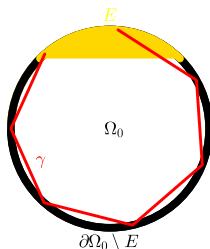
$$\int_0^L e^{-2\lambda t} (q_1 - q_2)^{\wedge}(2\lambda, \gamma(t)) dt = 0.$$

Here  $(\cdot)^{\wedge}$  is the Fourier transform in the  $x_1$  variable, and  $q_1 - q_2$  is extended by zero to  $\mathbb{R}^3 \setminus \overline{\Omega}$ .

# Beyond the convex hull

## Question

Let  $\Omega_0 \subset \mathbb{R}^n$  strictly convex and  $E \subset \partial\Omega_0$  open. Is a function  $f \in C(\overline{\Omega_0})$  determined by its integrals over broken rays starting and ending on  $E$ ?



- ▶ Eskin (2004): rays reflecting off convex obstacles
- ▶ Ilmavirta (2013): partial results for unit disk
- ▶ Hubenthal (2013): microlocal analysis for unit square
- ▶ related to (but not the same as) the v-line transform

# Components of proof

Need Carleman estimate with boundary terms:

$$\begin{aligned} & -\frac{1}{\tau} \int_{\partial\Omega} (\partial_\nu \varphi) e^{\pm 2\tau\varphi} |\partial_\nu v|^2 dS + \|e^{\pm\tau\varphi} v\|_{L^2(\Omega)}^2 \\ & \leq \frac{C}{\tau^2} \|e^{\pm\tau\varphi} (-\Delta + q)v\|_{L^2(\Omega)}^2, \quad v \in C^\infty(\bar{\Omega}), \quad v|_{\partial\Omega} = 0. \end{aligned}$$

Kenig-Sjöstrand-Uhlmann (2007) use convexified weights

$$\varphi_\varepsilon = \varphi + \frac{1}{\varepsilon\tau} \frac{\varphi^2}{2}, \quad \varepsilon > 0 \text{ small.}$$

Carleman estimate leads to solutions of  $(-\Delta + q)u = 0$  with

- ▶ good control on  $\{x \in \partial\Omega; \partial_\nu \varphi(x) < 0\}$
- ▶ *no control* on  $\{x \in \partial\Omega; \partial_\nu \varphi(x) = 0\}$ .



# Components of proof

Need Carleman estimate with boundary terms:

$$\begin{aligned} & -\frac{1}{\tau} \int_{\partial\Omega} (\partial_\nu \varphi) e^{\pm 2\tau\varphi} |\partial_\nu v|^2 dS + \|e^{\pm\tau\varphi} v\|_{L^2(\Omega)}^2 \\ & \leq \frac{C}{\tau^2} \|e^{\pm\tau\varphi} (-\Delta + q)v\|_{L^2(\Omega)}^2, \quad v \in C^\infty(\bar{\Omega}), \quad v|_{\partial\Omega} = 0. \end{aligned}$$

We use modified weights

$$\varphi_\varepsilon = \varphi + \frac{1}{\varepsilon\tau} \frac{\varphi^2}{2} + \frac{1}{\tau} \kappa, \quad \varepsilon > 0 \text{ small}, \quad \partial_\nu \kappa|_{\partial\Omega} < 0.$$

Carleman estimate leads to solutions of  $(-\Delta + q)u = 0$  with

- ▶ good control on  $\{x \in \partial\Omega; \partial_\nu \varphi(x) < 0\}$
- ▶ *weak control* on  $\{x \in \partial\Omega; \partial_\nu \varphi(x) = 0\}$ .

# Components of proof

Some arguments can also be done by reflection, e.g. if  $\Gamma_0$  is part of a graph

$$\Gamma_0 \subset \{(x_1, x_2, \eta(x_2)); x_1, x_2 \in \mathbb{R}\}$$

where  $\eta$  is a function  $\mathbb{R} \rightarrow \mathbb{R}$ . Flattening the boundary by  $x_3 \mapsto x_3 - \eta(x_2)$  transforms the Euclidean Laplacian into

$$\Delta_g \approx \sum_{j,k=1}^3 g^{jk} \partial_{x_j} \partial_{x_k}, \quad (g_{jk}(x)) = \begin{pmatrix} 1 & 0 \\ 0 & g_0(x_2, x_3) \end{pmatrix}.$$

Reflecting across  $x_3 = 0$  generates a *Lipschitz singularity* in the metric  $g_0$ . However, the singularity only appears in the lower right corner, and methods for the anisotropic Calderón problem (Kenig-S-Uhlmann 2011) still apply.

# Components of proof

Suppose  $\Omega$  is part of a cylinder  $\mathbb{R} \times \Omega_0$  and

$$\Gamma_0 \subset \mathbb{R} \times (\partial\Omega_0 \setminus E)$$

where  $\Omega_0 \subset \mathbb{R}^2$  and  $E \subset \partial\Omega_0$ . Use complex geometrical optics solutions as  $\tau \rightarrow \infty$ ,

$$u(x_1, x') \approx e^{\pm\tau x_1} v_\tau(x')$$

where  $v_\tau(x')$  is a *reflected Gaussian beam quasimode* in  $\Omega_0$ , concentrating near a broken ray  $\gamma$  with endpoints on  $E$ :

$$\begin{aligned} \|(-\Delta - \tau^2)v_\tau\|_{L^2(\Omega_0)} &= O(\tau^{-K}), \quad \|v_\tau\|_{L^2(\partial\Omega_0 \setminus E)} = O(\tau^{-K}), \\ |v_\tau|^2 dx' &\rightarrow \delta_\gamma. \end{aligned}$$

Cf. Dos Santos-Kurylev-Lassas-S (upcoming).

# Summary

In the Calderón problem with partial data for  $n \geq 3$ :

- ▶ possible to ignore measurements on sets that are part of *cylindrical sets*, *conical sets*, or *surfaces of revolution*
- ▶ *local uniqueness* results that determine coefficients near the measurement set
- ▶ *global uniqueness* under certain size or concavity conditions, or if the broken ray transform is invertible

Survey with Kenig: "Recent progress in the Calderón problem with partial data" (2013).

# Open questions

## Question (Local data for $n \geq 3$ )

If  $\Omega \subset \mathbb{R}^n$ ,  $n \geq 3$ , if  $\Gamma$  is any open subset of  $\partial\Omega$ , and if  $q_1, q_2 \in L^\infty(\Omega)$ , show that  $C_{q_1}^{\Gamma, \Gamma} = C_{q_2}^{\Gamma, \Gamma}$  implies  $q_1 = q_2$ .

## Question (Data on disjoint sets for $n = 2$ )

If  $\Omega \subset \mathbb{R}^2$ , if  $\Gamma_D$  and  $\Gamma_N$  are disjoint open subsets of  $\partial\Omega$ , and if  $q_1, q_2 \in L^\infty(\Omega)$ , show that  $C_{q_1}^{\Gamma_D, \Gamma_N} = C_{q_2}^{\Gamma_D, \Gamma_N}$  implies  $q_1 = q_2$ .